

The Instanton -Torus Knot duality

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Based on works, Bulycheva, Nechaev , A.G. 1409..

Milekhin , A.G. 1412....,

Milekhin, Sopenko, A.G. 1506...

To appear

Plan of the talk

- Introduction. Where the question comes from?
- Condensate in 5D SUSY QED and QCD from the torus knot invariants
- Different representations of torus knot invariants
In 5D gauge theory.
- Torus knot invariants in the dual systems
- Conclusion

Examples of condensates. SUSY

- Holomorphy - «instantons» contribute to the condensate
- Longstanding puzzle concerning the gluino condensate in $N=1$ SYM(NSVZ). If one compact dimension- fractional instanton-monopoles contribute
- Exact result. Squark and monopole condensates vanishes at Argyres-Douglas point in $N=1$ SQCD- deconfinement (Yung, Vainshtein, A.G)

Old and new questions

- What is the microscopic picture behind the condensate formation? Examples- squark and gluino condensates in SQCD due to the zero modes in the instanton ensemble(NSVZ 83).Topological sector in the theory
- Chiral condensate in QCD. Quasizero modes in the instanton-antiinstanton ensemble.

New tools

- New invariants of knots. Khovanov homologies and superpolynomials which generalize Jones and HOMFLY polynomials
- Seiberg-Witten solution to $N=2$ SYM. Nekrasov partition sums. Explicit results for the instanton sums in the Omega-background
- Topological phases of matter. Classification via ground state degeneracy+ holonomy of Berry phase or via entanglement entropy

Knot invariants in gauge theories

- $J(q,K) = \langle W(K) \rangle$ in $SU(2)$ 3d Chern-Simons theory — Jones polynomial of knot K (Witten, 89). Can be generalized to all $SU(N)$ groups- HOMFLY polynomial $H(a,q,K)$, $a=q^N$
- Generalization to superpolynomial $P(a,q,t,K)$ (Dunfeld, Gukov, Rasmussen 04). Three gradings (a,q,t) in the Hilbert space
- All knot polynomials are particular indices
 $P(a,q,t,K) = \dim H_{\{ijk\}} a^i q^j t^k$ counts the multiplicities of the BPS states (Gukov, Schwartz, Vafa 04) in some theory

Torus knots can be drawn on the torus surface.
 $T(n,m)$ corresponds to two windings around cycles

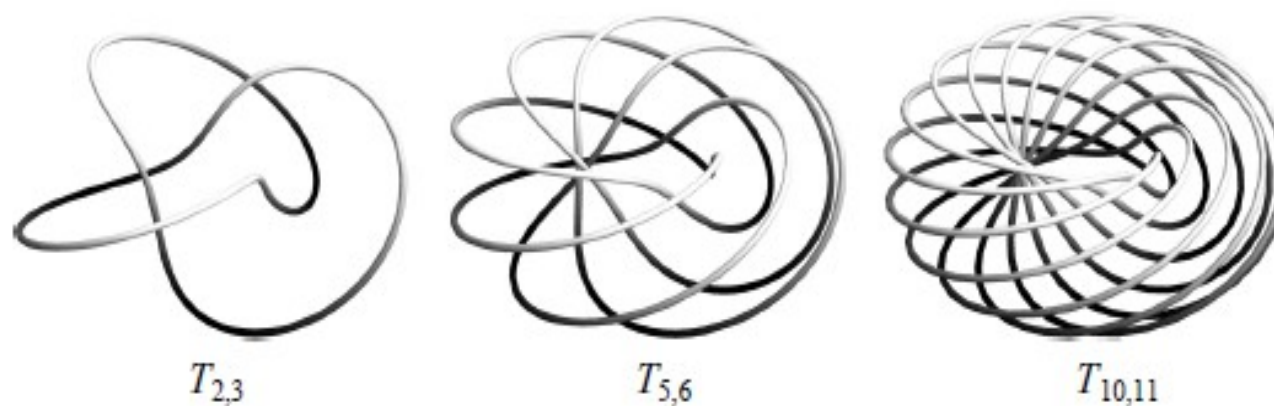


Figure 1: Few samples of torus knots from the series $T_{n,n+1}$:
 $T_{2,3}$, $T_{5,6}$, $T_{10,11}$.

Knot invariants. New approaches

- The old evaluation- it is vev of electric Wilson loop along the knot in 3d CS. Is there the **S-dual** «magnetic» version of the knot invariants? (Witten 10, Witten, Gaiotto 11). Witten's idea — knot invariants somehow count instantons in 4d and 5d $N=4$ SUSY gauge theory. Only partial success.
- The superpolynomials of the **torus knots** are expressed as very specific integrals over moduli space of points in C^2 (E.Gorsky-Negut , 13). Way to zero-size instantons in 5d theory.

5D SQED and SQCD

- Consider the $U(1)$ 5d SUSY gauge theory with $N_f=2$ or $N_f=3$. One dimension is compact S^1 . Add 5d CS term $k \text{ AFF}$. Introduce Omega-deformation = two independent rotations (angular velocities) in R^4 .
- There is the explicit answer for the instanton partition function in this theory due to Nekrasov localization
- Surprise. The condensate of the massless flavor « is sum over the invariants of the $T(n,m)$ torus knots in the momentum space »

Some facts on 5d SQED

- The BPS particles in the theory are W-bosons, instantons. Due to the CS term the instanton charge induces the electric charge
- Complicated dyonic instantons(both charges). Even more complicated states with 3 charges(+flavor). Not fully classified. Monopoles are loops(monopole particles lifted to 5d)
- One -loop effect of all BPS particles in 5d D with compact dimension reproduces all instanton partition sum in D=4 SYM theory(Nekrasov-Lawrence)

$$\Omega^m = \Omega^{mn} x_n, \quad \Omega^{mn} = \frac{1}{2\sqrt{2}} \begin{pmatrix} 0 & i\epsilon_1 & 0 & 0 \\ -i\epsilon_1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i\epsilon_2 \\ 0 & 0 & i\epsilon_2 & 0 \end{pmatrix}.$$

Omega-deformation by external field

$$\begin{aligned} \mathcal{L}_m = & -\frac{1}{4g^2} F_{mn} F^{mn} + \frac{1}{g^2} (\partial_m \phi + F_{mn} \Omega^n) (\partial^m \phi - F^{mn} \Omega_n) + \\ & \frac{1}{2} |D_m Q|^2 + \frac{1}{2} |D_m \tilde{Q}|^2 + \frac{2}{g^2} (i \partial_m (\Omega^m \bar{\phi} + \Omega^m \phi) + g^2 (\bar{Q} Q - \tilde{Q} \tilde{Q}))^2 + \\ & \frac{1}{2} |(\phi - m - i \Omega^m D_m) q|^2 + \frac{1}{2} |(\phi - \tilde{m} - i \Omega^m D_m) \tilde{q}|^2 + 2g^2 |\tilde{q} q|^2 \end{aligned}$$

Spectrum of BPS particles

$$Z = \frac{1}{g^2} n_I + n_e a + \sum_i n_{f_i} m_{f_i}$$

String web (Aharony, Hanany, Kol)

Instanton-torus knot duality,

$$\frac{e^{\beta M}}{(1+a)\beta^2} \frac{d^2 Z_{nek}(q, t, m, M, m_a, Q)}{dM dm} \Big|_{m \rightarrow 0, M \rightarrow \infty} = \sum_n Q^n (tq)^{n/2} P_{n, nk+1}(q, t, a) \quad (2)$$

where m_a, m, M are masses of three hypemultiplets and Q is the counting parameter for the instantons. The mapping between the parameters at the lhs and rhs goes as follows

$$t = \exp(-\beta\epsilon_1) \quad (3)$$

$$q = \exp(-\beta\epsilon_2) \quad (4)$$

$$a = -\exp(\beta m) \quad (5)$$

$$Q = \exp(-\beta/g^2) \quad (6)$$

Instanton-torus knot duality

The superpolynomial is the complicated product in terms of the Young tableou

$$P(A, q, t)_{nk+1, n} = \sum_{\lambda: |\lambda|=n} \frac{t^{(k+1)\sum l} q^{(k+1)\sum a} (1-t)(1-q) \prod^{0,0} (1 - Aq^{-a'} t^{-l'}) \prod^{0,0} (1 - q^{a'} t^{l'}) (\sum q^{a'} t^{l'})}{\prod (q^a - t^{l+1}) \prod (t^l - q^{a+1})}$$

In this formula there is only one independent index — n (instanton number)

New findings

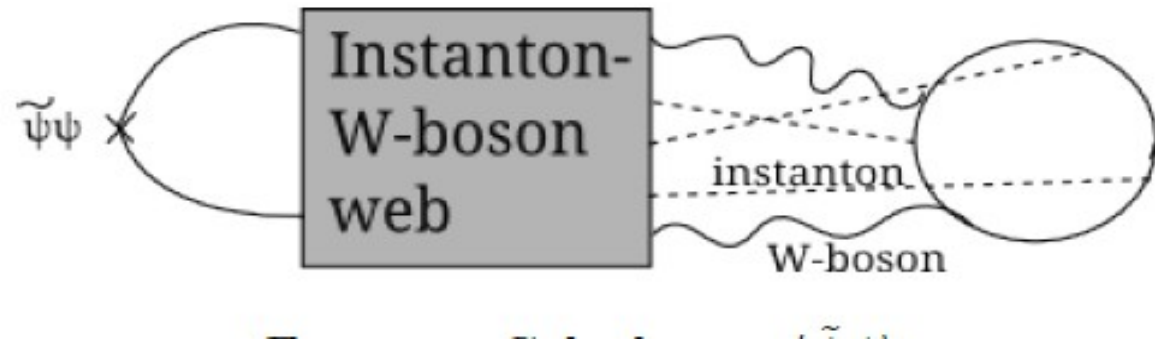
- The information about the knots is encoded in the condensate. Torus knots $T(m,n)$ are important. The physical identifications of the numbers: n - instanton charge, m -electric charge
- The physical variable is expressed in terms of the sum over the knots. The first example of such situation!
- The rank of the gauge group in the CS picture is the mass of the antifundamental(!)

Interpretation

- The knot invariants describe the multiplicity of states at fixed 4d quantum numbers (n,m)
- In some sense they count the 2d instantons on the nonabelian strings at fixed 4d instanton number. «Knotting the fermionic zero modes?»
- Instantons are membranes in the internal space and draw the knots on the «flavor branes». Knots live in the internal Calabi-Yau («momentum») space

Knot invariant as entropic factor

- The place of the knots in the diagrams



HOMFLY for generic (n,m) knots

- Consider the $N_f=2$ theory with Lagrangian brane. Count the contribution of states with (n,m) quantum numbers into condensate
- Consider the $SU(2)$ theory with $N_f=4$. Two masses fixed, one mass vanishes, one is arbitrary. Expand the condensate in series in two quantum numbers
- Consider $N_f=2$ $U(1)$ with fractional 5d CS number $k=m/n$. Extract n - instanton contribution

Sum over the (n,m)

- Double series for the condensate

$$\langle \tilde{\psi} \psi \rangle_{LB} = \left. \frac{\partial Z_{inst}}{\partial m_f} \right|_{m_f=0} = \sum_{n,m} Q_c^n z^m P_{n,nk+m}(A, q, t)$$

$$\sum_{\lambda: |\lambda|=n} \frac{t^{(k+1)\sum l} q^{(k+1)\sum a} (1-t) \prod^{0,0} (1 + Aq^{-a'} t^{-l'}) \prod^{0,0} (1 - q^{a'} t^{l'})}{\prod (q^a - t^{l+1}) \prod (t^l - q^{a+1})} \times$$

$$P(A, q, t)_{n,nk+m} =$$

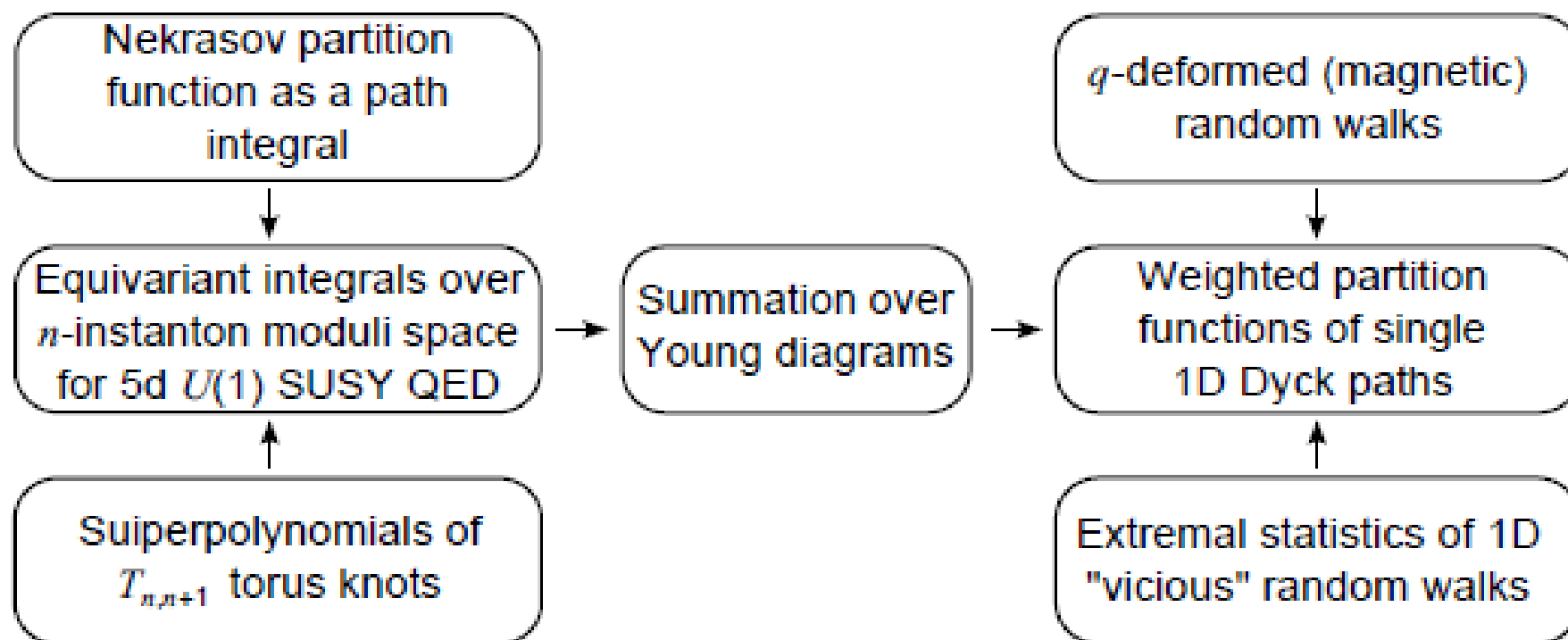
$$Coe f_{z^m} M(z)$$

where $M(z)$ is the contribution from the Lagrangian brane with zero framing:

$$M(z) = \prod_{j=1}^{l(\lambda)} \frac{1 - zt^{j-1} q^{\lambda_j}}{1 - zt^{j-1}}$$

Complimentary views on the instanton sums

Bulycheva, Nechaev, A.G. 1409



Knot invariants in the dual systems

- Nekrasov partition function in 5d is related to the q -Liouville conformal block. The derivative of the conformal block- generating function for knot invariants
- The Nekrasov partition function — wave function of the holomorphic Hamiltonian system. Perturbative and nonperturbative effects in SYM= similar effects in QM(example Toda, Basar-Dunne). Linking of pert and nonpert corrections via torus knot invariants (in progress)

Conclusion

- Just touch tip of the iceberg. Many surprises
- The «knotting» between electric degrees of freedom and instantons is important for the condensate formation
- Unexpected appearance of knot invariants in Liouville conformal blocks. General linking and knotting of perturbative and nonperturbative contributions.
- Some interplay with the solid state order parameters

A lot of open questions.....